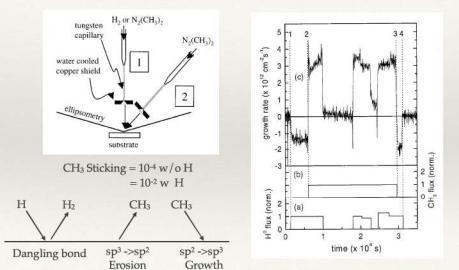
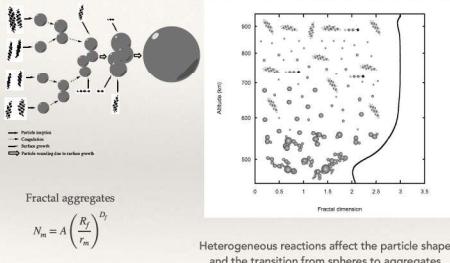


Simultaneous interaction of methyl radicals and atomic hydrogen with amorphous hydrogenated carbon films

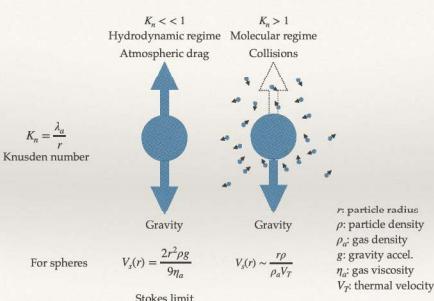
A. von Keudell,^a T. Schwarz-Selinger, and W. Jacob
Max-Planck-Institut für Plasmaphysik, EURATOM Association, Boltzmannstr. 2, 85748 Garching, Germany



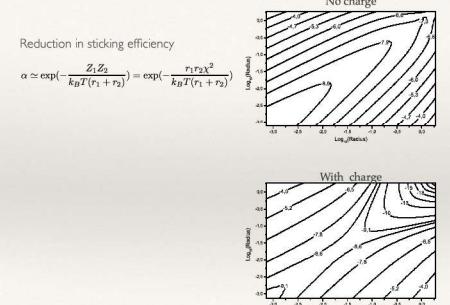
HETEROGENEOUS CHEMISTRY



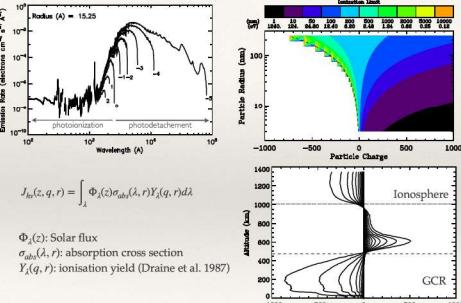
Settling velocity



Charge effects



Charge effects



Solution

Continuity equation for particles

$$\frac{\partial n_p(\nu_p)}{\partial t} = -\frac{1}{r^2} \frac{\partial(r^2 \Phi(\nu_p))}{\partial r} + \frac{1}{2} \int_{r_p}^{\infty} K(u, \nu_p - u) du - n(\nu_p) \int_0^{\infty} K(u, \nu_p) n(u) du + \frac{\partial n_p}{\partial t}_{\text{prod}}$$

Bin method: Solve for density of particles over a geometrically expanding volume grid.

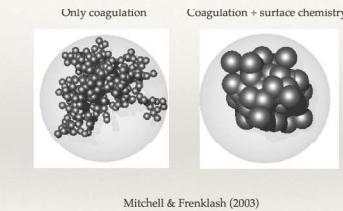
Slow but more informative.

Moments method: Assume a distribution form (e.g. lognormal) and solve for the moments of the distribution.

Requires assumptions, but much faster and integrable to GCMs!
(de Batz de Trenquelléon et al. 2025)

How do these surface processes affect the particle shape?

From combustion chemistry



Pure microphysics

Continuity equation for particles

$$\frac{\partial n_p(\nu_p)}{\partial t} = -\frac{1}{r^2} \frac{\partial(r^2 \Phi(\nu_p))}{\partial r} + \frac{1}{2} \int_{r_p}^{\infty} K(u, \nu_p - u) du - n(\nu_p) \int_0^{\infty} K(u, \nu_p) n(u) du + \frac{\partial n_p}{\partial t}_{\text{prod}}$$

Flux divergence

Growth from smaller particles

Loss to larger particles

Inception

n_p : number density particles with volume ν_p

$\Phi(\nu_p)$: flux of particles

$K(u, u')$: coagulation Kernel of particles with volumes u and u'

$$\Phi(\nu_p) = -V_s(\nu_p)n(\nu_p) - K_{ZZ}n_u \frac{\partial(n(\nu_p)/n_u)}{\partial r} \quad V_s(\nu_p): \text{settling velocity}$$

$$K_{ZZ}: \text{atmospheric mixing}$$

$$n_u: \text{atmospheric density}$$

Coagulation Kernels

Brownian: random collisions between particles

$$K_B(r, s) = \alpha_4 \pi (D_p + D_s)(r + s)\beta \quad (\text{for spheres})$$

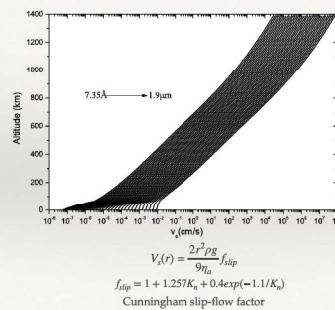
$$D_p = \frac{k_B T_{\text{slip}}}{6\pi\eta_a r} : \text{Particle diffusivity}$$

β, f : interpolation formulas between free-molecular ($K_n > 1$) and continuum regimes ($K_n < 1$)

α_4 : sticking efficiency

$$K_n = \frac{\lambda_0}{r} : \text{Knudsen number}$$

Settling velocity



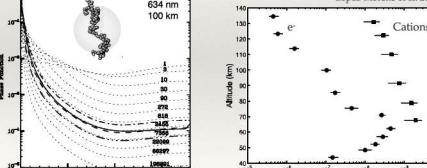
Charge effects

Reduction in sticking efficiency

$$\alpha \approx \exp(-\frac{Z_1 Z_2}{k_B T(r_1 + r_2)}) = \exp(-\frac{r_1 r_2 \chi^2}{k_B T(r_1 + r_2)})$$

$\chi = 15 \text{ e}^-/\mu\text{m}$

Tomasko et al. 2008, Lavvas et al. 2010



Reduction in sticking efficiency

$$\alpha \approx \exp(-\frac{Z_1 Z_2}{k_B T(r_1 + r_2)}) = \exp(-\frac{r_1 r_2 \chi^2}{k_B T(r_1 + r_2)})$$

Charge sources:

- collisions with ions & electrons
- photoionization
- photodetachment

$$J^+ = J_{pi} + J_{hv}$$

$$J^- = J_{ni} + J_e$$

Charge distributions

Particle Shape

Fractal aggregates

$$N_m = A \left(\frac{R_f}{r_m} \right)^D$$

$D_f = 1$ Linear

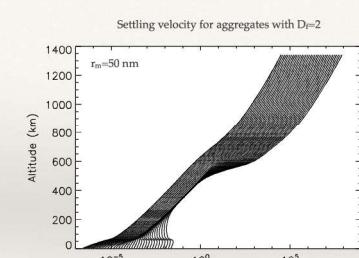
$D_f = 2$ Fluffy

$D_f = 3$ Compact



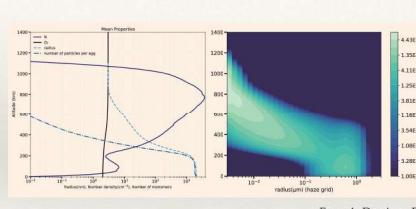
Growth of the cluster governed by:	$L \ll R$ (continuum regime)		$L \approx R$ (ballistic regime)	
	$D_s \approx 2.5$	$D_s \approx 3$	$D_s \approx 2.5$	$D_s \approx 2$
Particle-cluster aggregation (monomer-polymer)	$D_s \approx 1.75$			
Cluster-cluster aggregation (polymer-polymer)			$D_s \approx 2$	

$$r \rightarrow R_p = \frac{r^{3/D_f}}{r_m^{(3-D_f)/D_f}} \quad \text{for formulas based on spheres}$$



Particles settle with the velocity of the monomers!

Application to Titan



From A. Damien, PhD

Haze mass flux $\sim 3 \times 10^{-14} \text{ g cm}^{-2}\text{s}^{-1}$ corresponds to 30% of N₂ & CH₄ photolysis

Haze formation yield

Cloud formation